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On the Sharpness and Bias of Quantum Effects

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Abstract

The question of quantifying the *sharpness* (or unsharpness) of a quantum mechanical effect is investigated. Apart from sharpness, another property, *bias*, is found to be relevant for the joint measurability or coexistence of two effects. Measures of bias will be defined and examples given.

Dedication

The impossibility of measuring jointly certain pairs of observables is an intriguing non-classical feature of quantum theory that Pekka Lahti identified as a candidate for a rigorous formulation of the principle of complementarity. While he was investigating this fundamental no-go statement in the early 1980s, he learned from Peter Mittelstaedt that one of his students was aiming to prove the positive possibility of approximate joint measurements of complementary quantities such as position and momentum. Pekka joined our group as an Alexander von Humboldt Fellow, and together we found that a reconciliation between complementarity and (approximate) joint measurability is possible on the basis of the generalized representation of observables as positive operator measures (POMs). Since then we have pursued together our aspirations of understanding quantum mechanics and understanding Nature. I have benefited much from Pekka's intellectual rigor and have been privileged ever since to enjoy his warm humanity. It is a great pleasure to present this paper to Pekka as a token of thanks and friendship on the occasion of his 60th birthday, with all good wishes for many happy recurrences and productive years to come.

1 Introduction

The general description of quantum observables as positive operator measures (POMs) gives rise to a host of new operational possibilities not available within the set of standard observables (represented as projection valued measures). Here we focus on a particular issue of foundational significance: the possibility of joint measurements of certain pairs of noncommuting observables.

Two observables are considered to be jointly measurable if there is a third, *joint observable*, of which the given observables are marginals. According to a well known theorem of von Neumann [5], two observables represented as projection valued measures are jointly measurable if and only if they are mutually commuting [6]. Among pairs of general observables, commutativity is still a sufficient but no longer necessary condition for joint measurability.

Observables represented as projection valued measures are commonly understood to correspond to measurements with perfect accuracy; hence they can be called *sharp* observables. This characterization renders all other observables *unsharp*. If at least one of a pair of observables is sharp,

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then joint measurability cannot hold unless the two observables commute [3]. It follows that for two *noncommuting* observables to be jointly measurable, it is necessary that *both* of them are unsharp.

The problem thus arises of determining the factors that are relevant for the characterization of jointly measurable pairs of noncommuting observables. In light of the above general observation one would expect a trade-off to hold between the *degrees* of noncommutativity and sharpness within the set of jointly measurable pairs of observables. This leads to the task of defining appropriate measures of (un)sharpness. We will focus on the case of *simple observables*, that is, POMs generated by a resolution of the identity 1 of the form $\{A, 1 - A\}$, where both A and $1 - A$ are positive operators, also referred to as *effects*.

We will consider candidates of sharpness measures in 2-dimensional Hilbert spaces, to begin with. A recently found [2, 4, 7] criterion of the joint measurability of pairs of qubit effects is seen to involve a trade-off between the degrees of noncommutativity, sharpness and yet another quantity called bias. The task of extending measures of sharpness and bias to arbitrary Hilbert spaces is not entirely trivial due to ambiguity and the fact that the 2-dimensional case is too simple to reveal relevant features. Nevertheless we have been able to identify several distinct measures of sharpness and bias, applicable in Hilbert spaces of arbitrary dimensions.

2 Preliminaries

Our investigation is based on the usual quantum mechanical description of a physical system represented by a complex separable Hilbert space \mathcal{H} with inner product $\langle \varphi | \psi \rangle$, $\varphi, \psi \in \mathcal{H}$. States are represented by positive operators T of trace equal to one, the convex set of all states being denoted $\mathcal{S}(\mathcal{H})$. The extremal elements of $\mathcal{S}(\mathcal{H})$ are the vector states, that is the rank-one projections $T = |\varphi\rangle\langle\varphi|$, where φ is any unit vector of \mathcal{H} . An *effect* is a selfadjoint operator A on \mathcal{H} satisfying $0 \leq A \leq 1$. Here 0 and 1 are the null and identity operators, respectively, and the partial order $A \leq B$ is defined as $\langle \varphi | A \varphi \rangle \leq \langle \varphi | B \varphi \rangle$ for all $\varphi \in \mathcal{H}$. An effect A together with a state T gives the number $\text{tr}[TA] \in [0, 1]$, which is the probability for a measurement outcome represented by A to occur in a measurement performed on the state T .

The set of effects $\mathcal{E}(\mathcal{H})$ is thus the operator interval $[0, 1]$ with respect to \leq . $\mathcal{E}(\mathcal{H})$ is a convex set and its extremal elements are exactly the (orthogonal) projections ($A = A^2$). $\mathcal{E}(\mathcal{H})$ contains the convex subset of *trivial* effects $\lambda 1$, $\lambda \in [0, 1]$.

Projections will be called *sharp*, or *crisp effects*; all other effects are called *unsharp* or *fuzzy*. The unsharp trivial effects represent the extreme case of unsharpness; their expectation values provide no information about the state of a system. We note that an effect is a nontrivial sharp effect if and only if its spectrum consists of the two maximally separated elements 0,1; an effect is trivial if and only if its spectrum is a singleton. The intersection of the sets of sharp and trivial effects is $\{0, 1\}$.

The *complement* of an effect A is defined as $A' := 1 - A$. An algebraic relation that distinguishes sharp and unsharp effects arises from the inequality $A^2 \leq A$ which characterizes the effects among the selfadjoint operators. Hence we have $AA' \geq 0$ exactly when A is an effect, and an effect A is sharp exactly when $AA' = 0$.

The product of A and A' can be written as $AA' = A^{1/2}A'A^{1/2}$, and this suggests the following operational interpretation of the sharpness or unsharpness of an effect. Let ϕ_L^A be the Lüders operation associated with the effect A , corresponding to an ideal measurement (see [1]) of the simple observable given by A, A' , that is, $\phi_L^A(T) = A^{1/2}TA^{1/2}$ for any state T . The probability of a measurement giving an outcome associated with the effect A is given by $\text{tr}[TA] = \text{tr}[\phi_L^A(T)]$, and the sequential joint probability $\mathbf{p}_T(A, A')$ that a Lüders measurement of a simple observable \mathcal{A} repeated in immediate succession yields first the outcome associated with A and then the outcome associated with A' is

$$\mathbf{p}_T(A, A') = \text{tr}[\phi_L^A(T)A'] = \text{tr}[TA^{1/2}A'A^{1/2}] = \text{tr}[TAA']. \quad (1)$$

This joint probability is zero for all states if and only if A is a sharp effect. In this case, the measurement is repeatable for all states. For an unsharp effect the above joint probability is positive in some states. This joint probability becomes state-independent exactly when the effect A is trivial or has a spectrum of the form $\{a_0, (1 - a_0)\}$ for some $a_0 \in [0, 1]$. In the latter case A is of the form $A = a_0 P + (1 - a_0) P'$, where P is a projection and P' its complement.

3 Measures of sharpness and bias

3.1 Defining properties of a sharpness measure

We seek a definition of a measure of the *sharpness* of an effect A (and its associated simple observable $\mathcal{A} = \{A, A'\}$) which singles out and distinguishes the trivial effects on the one hand and the nontrivial sharp ones on the other hand. Specifically a function $\mathcal{E}(\mathcal{H}) \ni A \mapsto \mathfrak{S}(A)$ will be accepted as a sharpness measure if it satisfies (at least) the following requirements:

- (S1) $0 \leq \mathfrak{S}(A) \leq 1$;
- (S2) $\mathfrak{S}(A) = 0$ if and only if A is a trivial effect;
- (S3) $\mathfrak{S}(A) = 1$ if and only if A is a nontrivial projection;
- (S4) $\mathfrak{S}(A') = \mathfrak{S}(A)$;
- (S5) $\mathfrak{S}(CAC^{-1}) = \mathfrak{S}(A)$ for all invertible operators C ;
- (S6) $A \mapsto \mathfrak{S}(A)$ is (norm) continuous.

The function $A \mapsto \mathfrak{U}(A) = 1 - \mathfrak{S}(A)$ can then be taken as a measure of unsharpness or fuzziness. It satisfies a similar set of conditions, with appropriate adjustments.

Condition (S1) is merely a convention. (S2) and (S3) are the decisive properties of any sharpness measure once (S1) is adopted. (S4) and (S5) are motivated by the idea that the sharpness of an effect A is determined by the distribution of the spectrum σ_A of A . (S6) reflects the idea that a small change of A results in a small change of the degree of sharpness. (S5) and (S6) are automatically satisfied for sharpness measures $\mathfrak{S}(A)$ defined as functions of norm expressions such as $\|A\|$, $\|A'\|$, $\|AA'\|$.

We will see that the above list does not single out a unique sharpness measure. One could thus conceive of further properties that a sharpness measure may be required to have. For example, it could be argued that a convex combination of two effects cannot be sharper than these two effects themselves; this would be true if $A \mapsto \mathfrak{S}(A)$ is an affine functional. In view of the operational interpretation of unsharpness considered in the preceding section, we may have to be prepared to take into account that a measure $\mathfrak{S}(A)$ will depend on properties of AA' .

3.2 Bias of an effect

A yes-no experiment would be regarded as *biased* if one of the outcomes turned out to be preferred, whatever the preparation. We may thus define an effect A to be *unbiased* if for every state T there is a state T' such that $\text{tr}TA = \text{tr}T'A'$. This is equivalent to saying that the convex hulls of σ_A and $\sigma_{A'}$ are identical, or that the midpoints of σ_A and $\sigma_{A'}$ are the same. For later use we define the minimum, maximum, width, and midpoint of σ_A , respectively as:

$$M(\sigma_A) := \max(\sigma_A) = \|A\|; \quad (2)$$

$$m(\sigma_A) := \min(\sigma_A) = 1 - \max(\sigma_{A'}) = 1 - \|A'\|; \quad (3)$$

$$\mathfrak{W}(\sigma_A) := M(\sigma_A) - m(\sigma_A) = \|A\| + \|A'\| - 1; \quad (4)$$

$$\mu(\sigma_A) := \frac{1}{2} (M(\sigma_A) + m(\sigma_A)) = \frac{1}{2} (\|A\| + 1 - \|A'\|). \quad (5)$$

Then A is unbiased according to the above condition if $\mu(\sigma_A) = \frac{1}{2}$. We note that the set of unbiased effects is a convex subset of $\mathcal{E}(\mathcal{H})$. The *maximally biased* effects are 0 and 1.

A measure of bias $\mathcal{E}(\mathcal{H}) \ni A \mapsto \mathcal{B}(A)$ will be understood as a function that satisfies (at least) the following conditions:

- (B1) $-1 \leq \mathcal{B}(A) \leq 1$;
- (B2) $\mathcal{B}(A) = 0$ if and only if A is unbiased, i.e., $\mu(\sigma_A) = \frac{1}{2}$;
- (B3) $\mathcal{B}(A) = 1$ if and only if $A = 1$, and $\mathcal{B}(A) = -1$ if and only if $A = 0$;
- (B4) $\mathcal{B}(A) = -\mathcal{B}(A')$;
- (B5) $\mathcal{B}(CAC^{-1}) = \mathcal{B}(A)$ for all invertible operators C ;
- (B6) $A \mapsto \mathcal{B}(A)$ is (norm) continuous.

The motivation for these postulates is similar to those for sharpness measures. The last two can be secured by defining $\mathcal{B}(A)$ as a function of quantities such as $\|A\|$, $\|A'\|$, $\|AA'\|$.

4 Sharpness and bias measures for qubit effects

4.1 Examples of sharpness measures

We first construct examples of sharpness measures for qubit effects.

A qubit is described in a 2-dimensional Hilbert space, $\mathcal{H} = \mathbb{C}^2$. We use the Pauli basis $\{1, \sigma_1, \sigma_2, \sigma_3\}$ to represent an operator A as $A = a_0 1 + \mathbf{a} \cdot \boldsymbol{\sigma}$. Operator A is selfadjoint if and only if $(a_0, \mathbf{a}) \in \mathbb{R}^4$. A state operator T is given by $T = \frac{1}{2}(1 + \mathbf{t} \cdot \boldsymbol{\sigma})$, where the Euclidean norm of \mathbf{t} satisfies $0 \leq |\mathbf{t}| \leq 1$. A is an effect if and only if its eigenvalues are between 0 and 1, that is, $0 \leq a_0 \pm |\mathbf{a}| \leq 1$. An effect A is a projection if and only if $a_0 = |\mathbf{a}| = \frac{1}{2}$.

We recall that an effect A is nontrivial and sharp exactly when its spectrum $\sigma_A = \{0, 1\}$, and A is trivial exactly when $\sigma_A = \{a_0\}$. We seek a function $\mathcal{E}(\mathbb{C}^2) \ni A \mapsto \mathfrak{S}(A) \in [0, 1]$ which assumes value 1 exactly in the former case and value 0 exactly in the latter case. Such a function is easily obtained: we simply take its value to be the *spectral width*, that is, the difference between the largest and smallest eigenvalue:

$$\mathfrak{S}_a^{(2)}(A) := \mathfrak{W}(\sigma_A) = (a_0 + |\mathbf{a}|) - (a_0 - |\mathbf{a}|) = 2|\mathbf{a}|. \quad (6)$$

The superscript indicates that this quantity is specific to the case of a 2-dimensional Hilbert space. It is easily verified that $\mathfrak{S}_a^{(2)}$ satisfies all conditions (S1)-(S6); in addition, it is convex.

This function extends in an obvious way to higher dimensions:

$$\mathfrak{S}_a(A) := \mathfrak{W}(\sigma_A) = \|A\| - (1 - \|A'\|) = \|A\| + \|A'\| - 1. \quad (7)$$

In $\mathcal{H} = \mathbb{C}^2$ we have $\mathfrak{S}_a = \mathfrak{S}_a^{(2)}$, and this function is a sharpness measure. However, in higher dimensions \mathfrak{S}_a fails to satisfy (S3): if A is an effect with spectrum $\{0, \alpha, 1\}$, with $0 < \alpha < 1$, then $\mathfrak{S}_a(A) = 1$ although A is not sharp.

Another simple function that constitutes a sharpness measure in \mathbb{C}^2 is the following:

$$\mathfrak{S}_b^{(2)}(A) := 4 \min(a_0, 1 - a_0) |\mathbf{a}| = 2|\mathbf{a}| [1 - |2a_0 - 1|]. \quad (8)$$

We note that $A = a_0 1 + \mathbf{a} \cdot \boldsymbol{\sigma}$ is an effect if and only if

$$0 \leq |\mathbf{a}| \leq \min(a_0, 1 - a_0). \quad (9)$$

Putting $\xi := \min(a_0, 1 - a_0)$, $\eta := |\mathbf{a}|$, it is straightforward to see that the function $(\xi, \eta) \mapsto 4\xi\eta$, defined on the domain

$$D := \{(\xi, \eta) : 0 \leq \xi \leq \frac{1}{2}, 0 \leq \eta \leq \xi\}, \quad (10)$$

assumes its minimum 0 exactly on the line segment in D given by $\eta = 0$, while the maximum 1 is reached exactly at the point $(\xi, \eta) = (\frac{1}{2}, \frac{1}{2})$. This proves the properties (S1), (S2) and (S3). The symmetry (S4) is given by construction. Since

$$a_0 = \frac{1}{2}(a_0 + |\mathbf{a}|) + \frac{1}{2}(a_0 - |\mathbf{a}|) = \mu(\sigma_A), \quad |\mathbf{a}| = \frac{1}{2}(a_0 + |\mathbf{a}|) - \frac{1}{2}(a_0 - |\mathbf{a}|) = \frac{1}{2}\mathfrak{W}(\sigma_A), \quad (11)$$

it is seen that $\mathfrak{S}_b^{(2)}$ depends only on the eigenvalues of A and thus (S5) is fulfilled. For the same reason, $\mathfrak{S}_2^{(2)}$ is norm continuous.

We can characterize a sharpness measure $\mathfrak{S}^{(2)}$ in \mathbb{C}^2 more systematically as follows. $\mathfrak{S}^{(2)}(A)$ should depend on the eigenvalues of A (in fulfillment of (S5)) and thus can be expressed as a function $f(a_0, |\mathbf{a}|)$ in view of (11). The domain of f is given by Eq. (9). Next, the condition (S4) reads $f(a_0, |\mathbf{a}|) = f(1 - a_0, |\mathbf{a}|)$, which entails that $\mathfrak{S}^{(2)}$ can in fact be written as a function of $\xi = \min(a_0, 1 - a_0)$ and $\eta = |\mathbf{a}|$, so that $\mathfrak{S}^{(2)}(A) = f(\xi, \eta)$, with the domain now being D as given in (10). Any continuous function f on D with the property that the maximum 1 is assumed exactly on $(\frac{1}{2}, \frac{1}{2})$ and the minimum 0 is assumed exactly on the points $(\xi, 0)$ gives rise to a sharpness measure.

Returning to the example of $\mathfrak{S}_b^{(2)}$, we note that this can be written in a form that lends itself to generalization to arbitrary Hilbert spaces:

$$\begin{aligned} \mathfrak{S}_b(A) &= 4 \min \left[\frac{1}{2}(\|A\| + 1 - \|A'\|), 1 - \frac{1}{2}(\|A\| + 1 - \|A'\|) \right] \frac{1}{2}(\|A\| + \|A'\| - 1) \\ &= \min [1 + \|A\| - \|A'\|, 1 - (\|A\| - \|A'\|)] (\|A\| + \|A'\| - 1) \\ &= 2 \min(\mu(\sigma_A), \mu(\sigma_{A'})) \mathfrak{W}(\sigma_A) \\ &= (1 - |\|A\| - \|A'\||) (\|A\| + \|A'\| - 1). \end{aligned} \quad (12)$$

As with \mathfrak{S}_a , we find again that $\mathfrak{S}_b(A) = 1$ if only $\|A\| = \|A'\| = 1$, that is, if $\{0, 1\} \subseteq \sigma_A$. So (S2) is violated if \mathcal{H} is at least 3-dimensional. We thus find that it seems less than straightforward to find a sharpness measure for Hilbert spaces of arbitrary dimensions by extending a measure suitable for \mathbb{C}^2 .

Moreover, a function like $\mathfrak{S}_b^{(2)}$ will not even have, in general, a unique extension. In fact, another extension of this function is given by the following:

$$\mathfrak{S}_1(A) := \|A\| + \|A'\| - [|AA'| + \|1 - AA'\|]. \quad (13)$$

The fact that in \mathbb{C}^2 we have $\mathfrak{S}_1 = \mathfrak{S}_b^{(2)}$ is easily verified by noting that

$$\|AA'\| = \frac{1}{4} - \min \left((\|A\| - \frac{1}{2})^2, (\|A'\| - \frac{1}{2})^2 \right), \quad \|1 - AA'\| = \frac{3}{4} + \max \left((\|A\| - \frac{1}{2})^2, (\|A'\| - \frac{1}{2})^2 \right) \quad (14)$$

We will show below that the function \mathfrak{S}_1 is in fact a sharpness measure.

4.2 Sharpness and bias measures from qubit coexistence

Another sharpness measure was found in connection with a criterion for the joint measurability for two qubit effects. Two effects A, B are jointly measurable, or coexistent, if there is a POM, called a joint observable for A, B , whose range contains A and B ; this ensures that both resolutions of identity $\{A, A'\}$ and $\{B, B'\}$ are contained in the range of the joint observable. This joint observable can always be taken to be generated by a resolution of the identity of the form $\{G_{11}, G_{12}, G_{21}, G_{22}\}$ so that

$$A = G_{11} + G_{12}, \quad B = G_{11} + G_{21}.$$

It was recently proven independently by three groups that effects $A = a_0 1 + \mathbf{a} \cdot \boldsymbol{\sigma}$ and $B = b_0 1 + \mathbf{b} \cdot \boldsymbol{\sigma}$ are coexistent exactly when a certain inequality holds [2, 4, 7]. This inequality can be cast in the form [2]

$$\frac{1}{2}[\mathcal{F}(2 - B) + \mathcal{B}(2 - \mathcal{F})] + (xy - 4\mathbf{a} \cdot \mathbf{b})^2 \geq 1. \quad (15)$$

Here the following abbreviations are used:

$$\mathcal{F} := \mathcal{F}^{(2)}(A)^2 + \mathcal{F}^{(2)}(B)^2; \quad (16)$$

$$\mathcal{B} := \mathcal{B}^{(2)}(A)^2 + \mathcal{B}^{(2)}(B)^2; \quad (17)$$

$$x := \mathcal{F}^{(2)}(A)\mathcal{B}^{(2)}(A) = 2a_0 - 1 \equiv \mathcal{B}_a^{(2)}(A); \quad (18)$$

$$y := \mathcal{F}^{(2)}(B)\mathcal{B}^{(2)}(B) = 2b_0 - 1 \equiv \mathcal{B}_a^{(2)}(B); \quad (19)$$

$$\mathcal{F}^{(2)}(A) := \sqrt{a_0^2 - |\mathbf{a}|^2} + \sqrt{(1 - a_0)^2 - |\mathbf{a}|^2}; \quad (20)$$

$$\mathcal{B}^{(2)}(A) := \sqrt{a_0^2 - |\mathbf{a}|^2} - \sqrt{(1 - a_0)^2 - |\mathbf{a}|^2}. \quad (21)$$

$\mathcal{F}^{(2)}(B)$ and $\mathcal{B}^{(2)}(A)$ are defined similarly. It has been shown [4, 2] that the following is a sharpness measure:

$$\begin{aligned} \mathfrak{S}_c^{(2)}(A) &:= 1 - \mathcal{F}^{(2)}(A)^2 = 2 \left[a_0(1 - a_0) + |\mathbf{a}|^2 - \sqrt{(a_0^2 - |\mathbf{a}|^2)((1 - a_0)^2 - |\mathbf{a}|^2)} \right] \\ &= 2 \left[a_0(1 - a_0) + |\mathbf{a}|^2 \right] - \sqrt{4 \left[a_0(1 - a_0) + |\mathbf{a}|^2 \right]^2 - 4|\mathbf{a}|^2}. \end{aligned} \quad (22)$$

This function is even more complicated than the preceding ones. Still we will give a reformulation of this quantity so as to render it a sharpness measure for arbitrary Hilbert spaces.

It is also easy to see that $\mathcal{B}^{(2)}(A)$ and $\mathcal{B}_a^{(2)}(A)$ are both bias measures in \mathbb{C}^2 .

5 Sharpness and bias measures for arbitrary Hilbert spaces

We begin with a simple spectral characterization of an effect which arises as a generalization of inequality (9).

Proposition 1. *A selfadjoint bounded linear operator A in \mathcal{H} is an effect if and only if the following inequality holds:*

$$0 \leq \|A\| + \|1 - A\| - 1 \leq 1 - \left| \|A\| - \|1 - A\| \right|. \quad (23)$$

Note that for an effect A this relation can be written as (cf. Eq. 12)

$$0 \leq \frac{1}{2}\mathfrak{W}(\sigma_A) \leq \min(\mu(\sigma_A), \mu(\sigma_{A'})). \quad (24)$$

This shows that the statement is indeed quite obvious: this inequality ensures that the spectrum of A lies in the interval $[0, 1]$.

We proceed to construct some relatively simple measures of sharpness and bias for general Hilbert spaces.

Theorem 1. *The function \mathfrak{S}_0 given by*

$$\mathfrak{S}_0(A) := \|A\| \|A'\| - \|AA'\| \quad (25)$$

is a sharpness measure.

Proof. It is obvious that $\mathfrak{S}_0(A) \in [0, 1]$.

We write

$$M := \max \sigma_A = \|A\|, \quad m := \min \sigma_A = 1 - \|A'\|. \quad (26)$$

Now, if $\mathfrak{S}_0(A) = 0$, we have $\|A\| \|A'\| = \|AA'\|$. This reads:

$$M(1 - m) = \max_{\lambda \in \sigma_A} \lambda(1 - \lambda).$$

But here $\lambda \leq M$ and $1 - \lambda \leq 1 - m$, thus the above equation requires that $\lambda = M$ and $\lambda = m$, hence $m = M$ and $A = m1$. This proves (S2).

Next, suppose $\mathfrak{S}(A) = 1$. This entails $\|A\| \|A'\| = 1$ and $\|AA'\| = 0$, so that A is a nontrivial projection. Thus (S3) is verified. (S4), (S5) and (S6) are obvious. \square

Theorem 2. *The function*

$$\mathcal{E}(\mathcal{H}) \ni A \mapsto \mathcal{B}_0(A) := 2\mu(\sigma_A) - 1 = \|A\| - \|A'\| \quad (27)$$

is a bias measure.

The proof is trivial and will be omitted.

We can now give an equivalent way of phrasing the inequality (24) characterizing effects.

Corollary 1. *A bounded selfadjoint operator A is an effect if and only if*

$$\mathfrak{W}(\sigma_A) + |\mathcal{B}_0(A)| \leq 1. \quad (28)$$

Effects are thus distinguished by a trade-off between spectral width and bias.

In order to introduce the next sharpness measure, we note two relevant features of the spectrum σ_A of A . The first is the *spectral width*, which we denoted $\mathfrak{W}(\sigma_A)$ and defined as the length of the smallest interval containing σ_A . The second feature is the extent to which A deviates from the extreme cases of projections ($\sigma_A \subseteq \{0, 1\}$) or trivial effects ($\sigma_A = \{\lambda_0\}$). This latter feature will be called the *(spectral) dispersion* $\mathfrak{D}(\sigma_A)$ of A or of $\sigma_A \subseteq [0, 1]$; from the above discussion of the significance of the operator AA' , it is to be expected that $\mathfrak{D}(\sigma_A)$ is related to the spectral width of AA' . Thus we define:

$$\mathfrak{D}(\sigma_A) := \mathfrak{W}(\sigma_{AA'}). \quad (29)$$

To analyze these concepts, we recollect some simple observations.

$$\mathfrak{W}(\sigma_A) = \|A\| + \|A'\| - 1 = \mathfrak{W}(\sigma_{A'}); \quad (30)$$

$$\mathfrak{D}(\sigma_A) = \|AA'\| + \|1 - AA'\| - 1 = \mathfrak{D}(\sigma_{A'}). \quad (31)$$

We thus have $\mathfrak{S}_a(A) = \mathfrak{W}(A)$ and $\mathfrak{S}_1(A) = \mathfrak{W}(\sigma_A) - \mathfrak{D}(\sigma_A)$.

Lemma 1. *Let A be an effect. Then:*

$$\mathfrak{W}(\sigma_A) = 2 \min_{\kappa \in \mathbb{R}} \|A - \kappa 1\|. \quad (32)$$

The minimum is assumed at $\kappa = a_0$, where

$$a_0 = \frac{1}{2} [\|A\| - \|A'\| + 1] = \mu(\sigma_A). \quad (33)$$

Proof. We calculate $\|A - \kappa 1\| = \sup\{|\langle \varphi | A \varphi \rangle - \kappa| : \varphi \in \mathcal{H}, \|\varphi\| = 1\}$ for the various possible cases.

(i) $\kappa \geq \|A\|$:

$$\|A - \kappa 1\| = \kappa - \min \sigma_A \geq \|A\| - [1 - \|A'\|] \quad (\text{attained at } \kappa = \|A\|);$$

(ii) $0 \leq \|A\| - \kappa \leq \kappa - [1 - \|A'\|]$:

$$\|A - \kappa 1\| = \kappa - \min \sigma_A \geq \frac{1}{2} [\|A\| - [1 - \|A'\|]] \quad (\text{attained at } \kappa = \frac{1}{2} [\|A\| + [1 - \|A'\|]]);$$

(iii) $\|A\| - \kappa \geq \kappa - [1 - \|A'\|] \geq 0$:

$$\|A - \kappa 1\| = \|A\| - \kappa \geq \frac{1}{2} [\|A\| - [1 - \|A'\|]] \quad (\text{attained at } \kappa = \frac{1}{2} [\|A\| + [1 - \|A'\|]]);$$

(iv) $[1 - \|A'\|] \geq \kappa$:

$$\|A - \kappa 1\| = \|A\| - \kappa \geq \|A\| - [1 - \|A'\|] \quad (\text{attained at } \kappa = [1 - \|A'\|]).$$

From this it is seen that $\min_{\kappa} \|A - \kappa 1\| = \frac{1}{2}(\|A\| + \|A'\| - 1) = \frac{1}{2}\mathfrak{W}(\sigma_A)$.

Using the value $\kappa = a_0$, we find

$$\|A - a_0 1\| = \max_{\lambda \in \sigma_A} |\lambda - a_0| = \|A\| - a_0 = \frac{1}{2}[\|A\| + \|A'\| - 1].$$

□

This shows that the spectral width of an effect A is twice the minimal norm distance of A from the set of trivial effects.

Proposition 2. *Let A be an effect; then:*

- (a) $0 \leq \mathfrak{W}(\sigma_A) \leq 1$;
- (b) $\mathfrak{W}(\sigma_A) = 1$ iff $\{0, 1\} \subseteq \sigma_A$;
- (c) $\mathfrak{W}(\sigma_A) = 0$ iff A is trivial.

Proof. Inequality (a) is an immediate consequence of the definition of $\mathfrak{W}(\sigma_A)$.

(b): For $\mathfrak{W}(\sigma_A)$ to be equal to 1 it is necessary and sufficient that $\max \sigma_A = 1$ and $\min \sigma_A = 0$.

(c): $\mathfrak{W}(\sigma_A) = 0$ is equivalent to $\max \sigma_A = \min \sigma_A$, which means that σ_A is a one-point set or that A is a constant operator. □

Lemma 2. *For an effect A the operator AA' has the following properties:*

- (a) $0 \leq AA' \leq \frac{1}{4}1$;
- (b) $\frac{1}{4} - \max\{(\|A\| - \frac{1}{2})^2, (\|A'\| - \frac{1}{2})^2\} \leq \|AA'\| \leq \frac{1}{4}$;
- (c) $\|AA'\| = \frac{1}{4}$ if and only if $\frac{1}{2} \in \sigma_A$;
- (d) $\|1 - AA'\| = \frac{3}{4} + \max\{(\|A\| - \frac{1}{2})^2, (\|A'\| - \frac{1}{2})^2\}$.

Proof. Let $\lambda \mapsto E_{\lambda}^A$ denote the spectral family of A . Then

$$AA' = \int \lambda(1 - \lambda) dE_{\lambda}^A = \frac{1}{4}1 - \int (\lambda - \frac{1}{2})^2 dE_{\lambda}^A = \frac{1}{4}1 - (A - \frac{1}{2}1)^2. \quad (34)$$

Using this and the fact that $\lambda(1 - \lambda) \leq \frac{1}{4}$, with the maximum attained at $\lambda = \frac{1}{2}$, we obtain immediately (a), (b) and the sufficiency part of (c). If $\frac{1}{2} \notin \sigma_A$ then there is a $\lambda_0 \in \sigma_A$ which assumes the shortest positive distance of the closed set σ_A from $\frac{1}{2}$; it follows that $\|AA'\| = \lambda_0(1 - \lambda_0) < \frac{1}{4}$. This proves the necessity part of (c). To show (d), we note that $\lambda \mapsto \lambda(1 - \lambda)$ is increasing for $0 \leq \lambda \leq \frac{1}{2}$ and decreasing for $\frac{1}{2} \leq \lambda \leq 1$; then $\min\{\lambda(1 - \lambda) : \lambda \in \sigma_A\}$ is attained either at $\lambda = \min \sigma_A$ or $\lambda = \max \sigma_A$, hence:

$$\begin{aligned} \|1 - AA'\| &= \max\{1 - \|A\|(1 - \|A\|), 1 - \|A'\|(1 - \|A'\|)\} \\ &= \frac{3}{4} + \max\left\{\left(\|A\| - \frac{1}{2}\right)^2, \left(\|A'\| - \frac{1}{2}\right)^2\right\}. \end{aligned}$$

□

Proposition 3. *Let A be an effect. Then:*

(a) $\mathfrak{D}(\sigma_A)$ is given by

$$\mathfrak{D}(\sigma_A) = \max\left\{\left(\|A\| - \frac{1}{2}\right)^2, \left(\|A'\| - \frac{1}{2}\right)^2\right\} - \left(\frac{1}{4} - \|AA'\|\right); \quad (35)$$

- (b) $0 \leq \mathfrak{D}(\sigma_A) \leq \frac{1}{4}$;
- (c) $\mathfrak{D}(\sigma_A) = \frac{1}{4}$ iff $\frac{1}{2} \in \sigma_A$ and $1 \in \sigma_A$ or $0 \in \sigma_A$.
- (d) $\mathfrak{D}(\sigma_A) = 0$ iff $\sigma_A = \{\lambda\}$ or $\sigma_A = \{\lambda, 1 - \lambda\}$ (with $0 \leq \lambda \leq 1$).

Proof. The first equation is a direct consequence of Lemma 2.

It follows directly from the definition that $\mathfrak{D}(\sigma_A) \geq 0$. The maximum of $\mathfrak{D}(\sigma_A)$ is obtained by maximizing the positive term in Eq. (35) and minimizing the negative term: this happens when either $\|A\| = 1$ or $\|A'\| = 1$, and $\frac{1}{2} \in \sigma_A$. This yields $\max_{A \in \mathcal{E}(\mathcal{H})} \mathfrak{D}(\sigma_A) = \frac{1}{4}$.

Next, $\mathfrak{D}(\sigma_A) = 0$ means that AA' is a trivial effect (Lemma 1). Using again the spectral representation (34) of AA' , we see that AA' is a multiple of 1 if and only if σ_A is such that $\lambda(1-\lambda)$ is a constant $\frac{1}{4} - \epsilon^2$. This is the case if and only if either $\sigma_A = \{\lambda\}$ or $\sigma_A = \{\lambda, 1-\lambda\}$ (where $0 \leq \lambda \leq 1$). \square

This result shows that $\mathfrak{D}(\sigma_A)$ becomes large only if σ_A extends to at least one of the end points of the interval $[0,1]$ and to its center.

Theorem 3. *For any effect A the following holds true:*

- (a) $0 \leq \mathfrak{W}(\sigma_A) - \mathfrak{D}(\sigma_A) \leq 1$;
- (b) $\mathfrak{W}(\sigma_A) - \mathfrak{D}(\sigma_A) = 0$ if and only if A is trivial;
- (c) $\mathfrak{W}(\sigma_A) - \mathfrak{D}(\sigma_A) = 1$ if and only if A is a nontrivial projection.

Proof. We know already that $\mathfrak{W}(\sigma_A) \leq 1$ and $\mathfrak{D}(\sigma_A) \geq 0$; this gives $\mathfrak{W}(\sigma_A) - \mathfrak{D}(\sigma_A) \leq 1$.

To prove that this quantity is nonnegative we use the explicit form:

$$\begin{aligned} \mathfrak{W}(\sigma_A) - \mathfrak{D}(\sigma_A) &= \|A\| - \frac{1}{2} + \|A'\| - \frac{1}{2} + \left[\frac{1}{4} - \|AA'\|\right] \\ &\quad - \max\{(\|A\| - \frac{1}{2})^2, (\|A'\| - \frac{1}{2})^2\}. \end{aligned} \quad (36)$$

We consider the following cases.

- (i) $\epsilon \equiv \|A\| - \frac{1}{2} \geq 0$, $\epsilon' \equiv \|A'\| - \frac{1}{2} \geq 0$.

This case entails that $\sigma_A \subseteq [\frac{1}{2} - \epsilon', \frac{1}{2} + \epsilon]$, $\sigma_{A'} \subseteq [\frac{1}{2} - \epsilon, \frac{1}{2} + \epsilon']$, where the interval boundaries are the minima and maxima of the spectra. Then using $\|AA'\| \leq \frac{1}{4}$ we estimate:

$$\begin{aligned} \mathfrak{W}(\sigma_A) - \mathfrak{D}(\sigma_A) &= \epsilon + \epsilon' - \max\{\epsilon^2, \epsilon'^2\} + \left[\frac{1}{4} - \|AA'\|\right] \\ &\geq (\epsilon + \epsilon') - \frac{1}{2}[\epsilon^2 + \epsilon'^2 + |\epsilon^2 - \epsilon'^2|] \\ &= \begin{cases} \epsilon(1 - \epsilon) + \epsilon' & \text{if } \epsilon \geq \epsilon'; \\ \epsilon'(1 - \epsilon') + \epsilon & \text{if } \epsilon \leq \epsilon'. \end{cases} \end{aligned}$$

This becomes equal to 0 exactly for $\epsilon = \epsilon' = 0$, that is, $A = A' = \frac{1}{2}1$.

- (ii) $-\frac{1}{2} \leq -\epsilon \equiv \|A\| - \frac{1}{2} \leq 0 \leq \epsilon' \equiv \|A'\| - \frac{1}{2} \leq \frac{1}{2}$.

This condition entails that $\sigma_A \subseteq [\frac{1}{2} - \epsilon', \frac{1}{2} - \epsilon]$ and $\sigma_{A'} \subseteq [\frac{1}{2} + \epsilon, \frac{1}{2} + \epsilon']$. Here the bounds are the minima and maxima of the spectra. Furthermore, $\|AA'\| = \left\| \int \lambda(1-\lambda) dE_\lambda^A \right\| = \frac{1}{4} - \epsilon^2$. Then $\mathfrak{W}(\sigma_A) = \epsilon' - \epsilon \geq 0$, and $\mathfrak{D}(\sigma_A) = \max\{\epsilon^2, \epsilon'^2\} - (\frac{1}{4} - \|AA'\|) = \epsilon'^2 - \epsilon^2 = (\epsilon' - \epsilon)(\epsilon' + \epsilon)$. Thus,

$$\mathfrak{W}(\sigma_A) - \mathfrak{D}(\sigma_A) = (\epsilon' - \epsilon)[1 - (\epsilon + \epsilon')] \geq 0.$$

The latter expression is 0 exactly when $\epsilon' = \epsilon$, that is, when $\sigma_A = \{\frac{1}{2} - \epsilon\}$, i.e., when A is trivial.

- (iii) $-\frac{1}{2} \leq -\epsilon' := \|A'\| - \frac{1}{2} \leq 0 \leq \epsilon = \|A\| - \frac{1}{2} \leq \frac{1}{2}$.

This case is analogous to the previous one, with A and A' exchanged.

- (iv) $\|A\| < \frac{1}{2}$ and $\|A'\| < \frac{1}{2}$.

This case does not arise since $\|A\| < \frac{1}{2}$ implies $\|A'\| > \frac{1}{2}$.

Finally, assume that $\mathfrak{W}(\sigma_A) - \mathfrak{D}(\sigma_A) = 1$. This is equivalent to $\mathfrak{W}(\sigma_A) = 1$ and $\mathfrak{D}(\sigma_A) = 0$. From our previous results this is equivalent to $\{0, 1\} \subseteq \sigma_A$ and A being either trivial or having $\sigma_A = \{\frac{1}{2} \pm \epsilon\}$. The first condition rules out triviality and requires $\epsilon = 0$. It follows that A must be a nontrivial projection. Conversely, it is obvious that this condition entails that $\mathfrak{W}(\sigma_A) = 1$ and $\mathfrak{D}(\sigma_A) = 0$. \square

The following is now immediate.

Corollary 2. The function $\mathcal{E}(\mathcal{H}) \ni A \mapsto \mathfrak{S}_1(A)$, defined by

$$\mathfrak{S}_1(A) := \mathfrak{W}(\sigma_A) - \mathfrak{D}(\sigma_A) = \mathfrak{W}(\sigma_A) - \mathfrak{W}(\sigma_{AA'}), \quad (37)$$

satisfies requirements (S1)-(S6).

This proves that the extension \mathfrak{S}_1 of $\mathfrak{S}_b^{(2)}$ is a sharpness measure for Hilbert spaces of arbitrary dimension.

6 Further measures of sharpness and bias from the coexistence condition

There are various extension of the sharpness measure $\mathfrak{S}_c^{(2)}$. The first, \mathfrak{S}_2 , arises from the observation that in \mathbb{C}^2 the following equations hold:

$$\begin{aligned} 2\|A\| \|A'\| - \mathfrak{W}(\sigma_A) &= 2[a_0(1 - a_0) + |\mathbf{a}|^2]; \\ 2\mathfrak{S}_0(A) - \mathfrak{S}_1(A) &= 4|\mathbf{a}|^2. \end{aligned}$$

This leads us to define:

$$\mathfrak{S}_2(A) := 2\|A\| \|A'\| - \mathfrak{W}(\sigma_A) - \sqrt{[2\|A\| \|A'\| - \mathfrak{W}(\sigma_A)]^2 - [2\mathfrak{S}_0(A) - \mathfrak{S}_1(A)]}. \quad (38)$$

Theorem 4. The function $A \mapsto \mathfrak{S}_2(A)$ is a sharpness measure.

Proof. It suffices to prove (S1)-(S3). Write $\mathfrak{S}_2(A)$ in the form $\mathfrak{S}_2(A) = X - \sqrt{X^2 - Y}$. We first show that

$$0 \leq X \equiv 2\|A\| \|A'\| - \mathfrak{W}(\sigma_A) \leq 1, \quad 0 \leq Y \equiv 2\mathfrak{S}_0(A) - \mathfrak{S}_1(A) \leq X^2.$$

This ensures that $\mathfrak{S}_2(A) \in [0, 1]$. We observe:

$$X = 2\|A\| \|A'\| - \|A\| + (1 - \|A'\|) = 2M(1 - m) - M + m = M + m - Mm = M(1 - m) + m(1 - M) \in [0, 1].$$

Next we compute:

$$\begin{aligned} Y &= 2\mathfrak{S}_0(A) - \mathfrak{S}_1(A) = 2\|A\| \|A'\| - 2\|AA'\| - \mathfrak{W}(\sigma_A) + \mathfrak{W}(\sigma_{AA'}) \\ &= (2\|A\| \|A'\| - \mathfrak{W}(\sigma_A)) + (\mathfrak{W}(\sigma_{AA'}) - 2\|AA'\|) \\ &= X + \|1 - AA'\| - 1 - \|AA'\| = X - 2\mu(\sigma_{AA'}) \\ &= M(1 - m) + m(1 - M) + \frac{3}{4} + \max\left\{(M - \frac{1}{2})^2, (m - \frac{1}{2})^2\right\} - 1 - \frac{1}{4} + \min_{\lambda \in \sigma_A}(\lambda - \frac{1}{2})^2 \\ &= M(1 - m) + m(1 - M) - \frac{1}{2} + \min_{\lambda \in \sigma_A}(\lambda - \frac{1}{2})^2 + \max\left\{(M - \frac{1}{2})^2, (m - \frac{1}{2})^2\right\}. \end{aligned}$$

We consider the following cases for

$$\Delta \equiv (M - \frac{1}{2})^2 - (m - \frac{1}{2})^2 = (M - m)(M + m - 1) :$$

(I) $\Delta = 0$, i.e., $M = m$ or $M = 1 - m$;

(II) $\Delta > 0$, i.e. $M > m$ and $M > 1 - m$;

(III) $\Delta < 0$, i.e. $M > m$ and $M < 1 - m$.

Also put $\min_{\lambda \in \sigma_A}(\lambda - \frac{1}{2})^2 \equiv \delta$.

Case (I) with $M = m$: then $\delta = (M - \frac{1}{2})^2$ and $Y = 2M(1 - M) - \frac{1}{2} + 2(M - \frac{1}{2})^2 = 0$.

Case (I) with $M = 1 - m$: then $Y = 3(\frac{1}{4} - M(1 - M)) + \delta \geq 0$.

Case (II): In this case we must have $M > \frac{1}{2}$ (since $m < M$ and $M + m > 1$). Then one obtains:

$$Y = (M - 2m + \frac{1}{2})(M - \frac{1}{2}) + \delta = \begin{cases} (M - m + \frac{1}{2} - m)(M - \frac{1}{2}) + \delta \geq 0 & \text{if } m \leq \frac{1}{2}; \\ (M - \frac{1}{2} - 2(m - \frac{1}{2}))(M - \frac{1}{2}) + (m - \frac{1}{2})^2 = (M - m)^2 & \text{if } m \geq \frac{1}{2}. \end{cases}$$

Case (III): In this case we must have $m < \frac{1}{2}$. One obtains:

$$Y = (2M - m - \frac{1}{2})(\frac{1}{2} - m) + \delta = \begin{cases} (M - \frac{1}{2} + M - m)(\frac{1}{2} - m) + \delta \geq 0 & \text{if } M \geq \frac{1}{2}; \\ (\frac{1}{2} - m - 2(\frac{1}{2} - M))(\frac{1}{2} - m) + (\frac{1}{2} - M)^2 = (M - m)^2 & \text{if } M \leq \frac{1}{2}. \end{cases}$$

Thus we see that $Y \geq 0$ in all cases.

Next we show that $Y \leq X^2$. Note that $X = 2M(1 - m) - M + m = M - 2Mm + m$. Then

$$X^2 - Y = (M - 2Mm + m)^2 - Y.$$

We consider again the above cases, using expressions for Y obtained there.

Case (I) with $M = m$: here $Y = 0$, so $X^2 - Y = X^2 \geq 0$.

Case (I) with $M = 1 - m$: here we have

$$\begin{aligned} X^2 - Y &= (1 - 2M(1 - M))^2 - 3(\frac{1}{4} - M(1 - M)) - \delta \\ &= 4M^2(1 - M)^2 + (M - \frac{1}{2})^2 - \delta \geq 4M^2(1 - M)^2 \geq 0. \end{aligned}$$

Case (II): Here we obtain:

$$\begin{aligned} X^2 - Y &= (M - 2Mm + m)^2 - [(M - 2m + \frac{1}{2})(M - \frac{1}{2}) + \delta] \\ &= (M - 2Mm + m)^2 - (M - m)^2 + (m - \frac{1}{2})^2 - \delta \\ &= 2M(1 - m)2m(1 - M) + (m - \frac{1}{2})^2 - \min_{\lambda \in \sigma_A} (\lambda - \frac{1}{2})^2 \geq 0. \end{aligned}$$

Case (III): We have:

$$\begin{aligned} X^2 - Y &= (M - 2Mm + m)^2 - [(2M - m - \frac{1}{2})(\frac{1}{2} - m) + \delta] \\ &= (M - 2Mm + m)^2 - (M - m)^2 + (M - \frac{1}{2})^2 - \min_{\lambda \in \sigma_A} (\lambda - \frac{1}{2})^2 \geq 0. \end{aligned}$$

Finally we check for which effects A one has $\mathfrak{S}_2(A)$ equal to 0 or 1. First consider $\mathfrak{S}_2(A) = 0$. This implies that $Y = 2\mathfrak{S}_0(A) - \mathfrak{S}_1(A) = 0$. In case (I) this holds whenever $M = m$, and where $M = 1 - m$, it requires $M = m = \frac{1}{2}$. It is also easily verified that in the cases (II) and (III) we have $Y = 0$ exacty when $M = m$. Hence $\mathfrak{S}_2(A) = 0$ is equivalent to A being trivial.

Assume now that $\mathfrak{S}_2(A) = 1$. This is equivalent to $X = 1$ and $X^2 - Y = 0$. The first condition reads: $M - 2Mm + m = M(1 - m) + m(1 - M) = 1$, which can only be satisfied by $M = 1$, $m = 0$. Hence $\sigma_A \supseteq \{0, 1\}$. This only leaves case (I) with $M = 1 - m = 1$ and thus $X^2 - Y = 4M^2(1 - M)^2 + (M - \frac{1}{2})^2 - \delta = \frac{1}{4} - \delta = 0$. This finally entails $|\lambda - \frac{1}{2}| = \frac{1}{2}$ for all $\lambda \in \sigma_A$ and so $\sigma_A = \{0, 1\}$, that is, A is a nontrivial projection. \square

It would seem natural to try and define a measure \mathcal{B}_2 of bias associated with \mathfrak{S}_2 . Considering equations (20), (21) and (38), one would expect the following to define \mathcal{B}_2 :

$$1 - \mathcal{B}_2(A)^2 := 2\|A\| \|A'\| - \mathfrak{W}(\sigma_A) + \sqrt{[2\|A\| \|A'\| - \mathfrak{W}(\sigma_A)]^2 - [2\mathfrak{S}_0(A) - \mathfrak{S}_1(A)]}. \quad (39)$$

Using the notation introduced in the proof of theorem 4, we then have:

$$\mathfrak{G}_2(A) = X - \sqrt{X^2 - Y}, \quad 1 - \mathcal{B}_2(A) = X + \sqrt{X^2 - Y}. \quad (40)$$

Mimicking the definition of equation (18), we would expect the quantity x there to be given by

$$\begin{aligned} \mathcal{B}_1(A)^2 \equiv x^2 &= \mathcal{F}_2(A)^2 \mathcal{B}_2(A)^2 = (1 - X + \sqrt{X^2 - Y})(1 - X - \sqrt{X^2 - Y}) \\ &= (1 - X)^2 - (X^2 - Y) = 1 - 2X + Y. \end{aligned} \quad (41)$$

It is evident that $\mathcal{B}_1(A)^2 \geq 0$ is equivalent to $\mathcal{B}_2(A)^2 \geq 0$, or $1 - \mathcal{B}_2(A)^2 \leq 1$. However, it turns out that these inequalities are not satisfied for all effects A . In fact, using the identity established in the proof of theorem 4,

$$Y = X - 2\mu(\sigma_{AA'}), \quad (42)$$

we find:

$$\begin{aligned} 1 - 2X + Y &= 1 - X - 2\mu(\sigma_{AA'}) \\ &= -2(M - \tfrac{1}{2})(\tfrac{1}{2} - m) + \min_{\lambda \in \sigma_A} (\lambda - \tfrac{1}{2})^2 + \max \{ (M - \tfrac{1}{2})^2, (m - \tfrac{1}{2})^2 \} \end{aligned}$$

In the case where $M = 1 - m$ (A unbiased) we find:

$$1 - 2X + Y = -(M - \tfrac{1}{2})^2 + \min_{\lambda \in \sigma_A} (\lambda - \tfrac{1}{2})^2 < 0$$

whenever $\sigma_A \neq \{m, M\}$. It follows that the above expressions for \mathcal{B}_1 and \mathcal{B}_2 fail to give bias measures.

We thus return to the equations (20) and (21) and try to find ways of writing these in a form applicable to arbitrary Hilbert spaces. Observe that

$$a_0 = \mu(\sigma_A) = 1 - \mu(\sigma_{A'}),$$

which we can take as a redefinition of a_0 . There are at least two ways of rewriting the quantity $a_0^2 - |\mathbf{a}|^2$:

$$a_0^2 - |\mathbf{a}|^2 = M(\sigma_A)m(\sigma_A) = Mm.$$

or

$$a_0^2 - |\mathbf{a}|^2 = a_0^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2 = \max_{\lambda \in \sigma_A} (2a_0\lambda - \lambda^2) = \|2\mu(\sigma_A)A - A^2\|.$$

and similarly

$$(1 - a_0)^2 - |\mathbf{a}|^2 = \|2\mu(\sigma_{A'})A' - A'^2\|.$$

In the first case we obtain:

$$\begin{aligned} \mathcal{F}_3(A) &:= \sqrt{M(\sigma_A)m(\sigma_A)} + \sqrt{M(\sigma_{A'})m(\sigma_{A'})} = \sqrt{Mm} + \sqrt{(1 - M)(1 - m)}; \\ \mathcal{B}_3(A) &:= \sqrt{M(\sigma_A)m(\sigma_A)} - \sqrt{M(\sigma_{A'})m(\sigma_{A'})} = \sqrt{Mm} - \sqrt{(1 - M)(1 - m)}. \end{aligned} \quad (43)$$

In the second case we have (noting that $\min_{\lambda \in \sigma_{A'}} (\lambda - a_0)^2 = \min_{\lambda \in \sigma_A} (\lambda - a_0)^2$):

$$\begin{aligned} \mathcal{F}_4(A) &:= \sqrt{a_0^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2} + \sqrt{(1 - a_0)^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2}; \\ \mathcal{B}_4(A) &:= \sqrt{a_0^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2} - \sqrt{(1 - a_0)^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2}. \end{aligned} \quad (44)$$

It is not hard to show that $\mathcal{F}_3(A) \in [0, 1]$, $\mathcal{F}_4(A) \in [0, 1]$, and $\mathcal{B}_3(A) \in [-1, 1]$, $\mathcal{B}_4(A) \in [-1, 1]$. In particular, $\mathcal{F}_3(A) \leq 1$ is a consequence of the Cauchy-Schwarz inequality:

$$(\sqrt{M}, \sqrt{1-M}) \cdot (\sqrt{m}, \sqrt{1-m}) \leq \sqrt{M+(1-M)}\sqrt{m+(1-m)} = 1,$$

and equality is reached exactly when $M = m$. Thus the condition $\mathcal{F}_3(A) = 1$ holds if and only if $A = M1$, that is, A is trivial.

The equation $\mathcal{F}_3(A) = 0$ is equivalent to $m = 0$, $M = 1$. This does not require A to be a nontrivial projection, so \mathcal{F}_3 fails to be an unsharpness measure.

The equation $\mathcal{F}_4(A) = 0$ is equivalent to $a_0 = \min_{\lambda \in \sigma_A} |\lambda - a_0| = 1 - a_0$. This, in turn, is equivalent to $a_0 = \frac{1}{2} = \min_{\lambda \in \sigma_A} |\lambda - a_0|$, that is, $\sigma_A = \{0, 1\}$. Thus A is a nontrivial projection.

However, $\mathcal{F}_4(A) = 1$ is equivalent to $\min_{\lambda \in \sigma_A} |\lambda - a_0| = 0$, that is, $a_0 \in \sigma_A$. This does not require A to be trivial. Hence \mathcal{F}_4 fails to be an unsharpness measure.

The equation $\mathcal{B}_3(A) = 0$ is equivalent to $Mm = (1-M)(1-m)$, that is, $M + m = 1$, or $\mu(\sigma_A) = \frac{1}{2}$. Thus (B2) is satisfied. Further, $\mathcal{B}_3(A) = 1$ holds if and only if $M = m = 1$ or $A = 1$; likewise, $\mathcal{B}_3(A) = -1$ if and only if $M = m = 0$, i.e., $A = 0$. So (B3) holds.

Next, $\mathcal{B}_4(A) = 0$ is equivalent to $a_0 = 1 - a_0$, that is, $a_0 = \frac{1}{2}$, so (B2) holds. Finally, $\mathcal{B}_4(A) = 1$ ($\mathcal{B}_4(A) = -1$) if and only if $a_0 = 1$ ($1 - a_0 = 1$), and so $A = 1$ ($A = 0$). Thus (B3) is fulfilled.

It is possible to combine the virtues of \mathcal{F}_3 and \mathcal{F}_4 to obtain an unsharpness measure and an associated bias measure.

Theorem 5. *The following defines an unsharpness measure on $\mathcal{E}(\mathcal{H})$:*

$$\begin{aligned} \mathcal{F}_5(A) := & \sqrt{\frac{1}{2}(a_0^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2) + \frac{1}{2}Mm} \\ & + \sqrt{\frac{1}{2}((1 - a_0)^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2) + \frac{1}{2}(1 - M)(1 - m)}. \end{aligned} \quad (45)$$

Similarly, the following defines a bias measure on $\mathcal{E}(\mathcal{H})$:

$$\begin{aligned} \mathcal{B}_5(A) := & \sqrt{\frac{1}{2}(a_0^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2) + \frac{1}{2}Mm} \\ & - \sqrt{\frac{1}{2}((1 - a_0)^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2) + \frac{1}{2}(1 - M)(1 - m)}. \end{aligned} \quad (46)$$

Here the abbreviations $a_0 = \mu(\sigma_A)$, $M = \max(\sigma_A)$, $m = \min(\sigma_A)$ are used.

Proof. We note that

$$\begin{aligned} a_0^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2 &= \frac{1}{4}(M + m)^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2 \\ &= \frac{1}{4}(M - m) - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2 + Mm \geq Mm, \end{aligned} \quad (47)$$

and similarly

$$(1 - a_0)^2 - \min_{\lambda \in \sigma_A} (\lambda - a_0)^2 \geq (1 - M)(1 - m). \quad (48)$$

Note that equality holds in both cases if and only if $\sigma_A = \{M, m\}$. This entails that

$$\mathcal{F}_3(A) \leq \mathcal{F}_5(A) \leq \mathcal{F}_4(A) \quad (49)$$

and $\mathcal{F}_5(A) \in [0, 1]$, $\mathcal{B}_5(A) \in [-1, 1]$. It is now a simple consequence of inequalities (47)-(49) that $\mathcal{F}_5(A) = 1$ if and only if $\mathcal{F}_3(A) = \mathcal{F}_4(A) = 1$, and $\mathcal{F}_5(A) = 0$ if and only if $\mathcal{F}_3(A) = \mathcal{F}_4(A) = 0$. This ensures that \mathcal{F}_5 is in fact a sharpness measure. It is equally straightforward to verify that \mathcal{B}_5 is a bias measure. \square

7 Conclusion

We have determined a variety of measures of the sharpness and bias of an effect. It was found that generalization of such measures from \mathbb{C}^2 to arbitrary Hilbert spaces \mathcal{H} is not unambiguous in that there are different extensions. Neither are such generalizations entirely trivial to construct: we encountered some suggestive candidates which nevertheless failed to possess the desired properties of sharpness or bias measures, indicating that the two-dimensional case may not yield all relevant features of such measures.

A recently proven criterion for the coexistence of two qubit effects suggests that in general besides unsharpness, the bias of an effect is a significant quantity. This criterion involves certain measures of unsharpness and bias in a perfectly symmetric way, with closely related definitions, and we have found an extension of both measures to a general Hilbert space. It is an open question whether these or similar measures are relevant for the coexistence of effects in Hilbert spaces of dimension greater than two. It would be interesting to investigate the coexistence of more than two effects and of pairs of general POMs and to find out whether necessary and sufficient conditions can still be cast in the form of inequalities involving unsharpness and bias and possibly other quantities.

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References

- [1] P. Busch, P.J. Lahti, and P. Mittelstaedt. *The Quantum Theory of Measurement*. Springer-Verlag, Berlin, second revised edition, 1996.
- [2] P. Busch and H. J. Schmidt. Coexistence of qubit effects. <http://arxiv.org/abs/0802.4167v3>.
- [3] P. Kruszcynski and W.M. de Muynck. Compatibility of observables represented by positive operator-valued measures. *Journal of Mathematical Physics*, 28:1761–1763, 1987.
- [4] P. Stano, D. Reitzner, and T. Heinosaari. Coexistence of qubit effects. *Physical Review A*, 78:012315, 2008.
- [5] J. von Neumann. *Mathematische Grundlagen der Quantenmechanik*. Springer, Berlin, 1932. English translation: *Mathematical Foundations of Quantum Mechanics*, Princeton University Press, Princeton, 1955.
- [6] K. Ylinen. On a theorem of Gudder on joint distributions of observables. In P. Lahti and P. Mittelstaedt, editors, *Symposium on the Foundations of Modern Physics*, pages 691–694, Singapore, 1985. World Scientific.
- [7] S. Yu, N. L. Liu, L. Li, and C. H. Oh. Joint measurement of two unsharp observables of a qubit. <http://arxiv.org/abs/0805.1538>.